

Moments of β -ensembles

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- Eigenvalue joint probability density function:

$$P(\vec{\lambda}) = \frac{1}{\mathcal{N}} \prod_{i=1}^N w(\lambda_i) \prod_{1 \leq j < k \leq N} |\lambda_k - \lambda_j|^\beta$$

- Gaussian - $w(\lambda) = \exp\left(-\frac{\lambda^2}{2}\right)$; $\lambda \in \mathbb{R}$
 - GOE : $\beta = 1$
 - GUE : $\beta = 2$
- Laguerre - $w(\lambda) = \lambda^\gamma \exp(-\lambda)$; $\lambda \geq 0$
- Jacobi - $w(\lambda) = \lambda^{\gamma_1} (1 - \lambda)^{\gamma_2}$; $-1 \leq \lambda \leq 1$
- $\beta > 0$ as a continuous parameter
 - Log-gas
 - Tridiagonal realisation
 - Calogero-Sutherland models

Moments

- Let

$$\langle f(\vec{\lambda}) \rangle = \int_{\mathcal{D}} f(\vec{\lambda}) P(\vec{\lambda}) d\lambda_1 \dots d\lambda_N$$

- The moments are:

$$\mu_k = \left\langle \sum_{j=1}^N \lambda_j^k \right\rangle$$

- Why are moments of β -ensembles interesting?

- The p.d.f. of the position of the maximum of certain log-correlated processes on the interval (including the logarithm of the GUE characteristic polynomial) can be studied through it's moments - which requires calculating the moments of the $J\beta E$ [Fyodorov/Le Doussal 2015].
- Genus expansions ($G\beta E$)

$$\frac{1}{N} \mu_{2k} = \sum_{g=0}^{[k/2]} \frac{c(g; k)}{N^{2g}}$$

The Result

- Our formula for the moments in terms of the averages of Jack polynomials is given by:

$$\mu_k = \sum_{\nu \vdash k} c_\nu \left\langle C_\nu^{(\alpha)}(\vec{\lambda}) \right\rangle$$

- The coefficients c_ν are the inverses of Jack characters.

Symmetric Functions

- Symmetric functions are invariant under a permutation of their variables, they are usually indexed by partitions of integers.

- The partitions of 3 ($\nu \vdash 3$) are (3), (2, 1) and (1, 1, 1).

- Young diagrams: e.g. (2,1) is represented by



- Power sum functions:

$$p_{(2,1)}(\lambda_1, \lambda_2, \lambda_3) = (\lambda_1^2 + \lambda_2^2 + \lambda_3^2) \cdot (\lambda_1 + \lambda_2 + \lambda_3)$$

- Monomials:

$$m_{(2,1)}(\lambda_1, \lambda_2, \lambda_3) = \lambda_1^2 \lambda_2 + \lambda_1^2 \lambda_3 + \lambda_2^2 \lambda_1 + \lambda_2^2 \lambda_3 + \lambda_3^2 \lambda_1 + \lambda_3^2 \lambda_2$$

Jack Polynomials

- Jack polynomials:

$$C_{\nu}^{(\alpha)}(\vec{\lambda}) \propto m_{\nu}(\vec{\lambda}) + \sum_{\sigma \preceq \nu} u_{\nu, \sigma}^{\alpha} m_{\sigma}(\vec{\lambda})$$

- $P(\vec{\lambda})$ appears as the ground-state wave-function squared of a Calogero-Sutherland quantum particle system where β is the inverse temperature.
- The higher energy eigenfunctions are given by:

$$P(\vec{\lambda}) \cdot [\text{linear combinations of Jack polynomials}] \quad \text{with } \alpha = \frac{2}{\beta}$$

- Jack polynomials are orthogonal w.r.t. $P(\vec{\lambda})$ - the average of a Jack polynomial w.r.t. $P(\vec{\lambda})$ is known for the Laguerre, Jacobi and Gaussian ensembles [Baker/Forrester 1997].

Jack Characters

- Jack polynomials and power sum functions are related by a linear transformation:

$$C_{\nu}^{(\alpha)}(\vec{\lambda}) = \frac{\alpha^{|\nu|} |\nu|!}{j_{\nu}} \sum_{\sigma \perp |\nu|} \theta_{\sigma}^{\nu}(\alpha) p_{\sigma} \quad (1)$$

- The coefficients $\theta_{\sigma}^{\nu}(\alpha)$ are called Jack characters.
- Jack polynomials with $\alpha = 1$ (i.e. $\beta = 2$) are the Schur functions and $\theta_{\sigma}^{\nu}(1)$ are the characters of the irreducible representations of the symmetric group.

Overview of Derivation

- Begin with the generating function for the power sum polynomials:

$$Q(t) = \sum_{k \geq 1} p_k(\vec{\lambda}) t^{k-1} = \sum_{i=1}^N \sum_{k \geq 1} \lambda_i^k t^{k-1} = \sum_{i=1}^N \frac{d}{dt} \log \left(\frac{1}{1 - \lambda_i t} \right)$$

- Integrate:

$$\sum_{k \geq 1} \frac{p_k(\vec{\lambda})}{k} t^k = - \sum_{i=1}^N \log \left(\frac{1}{1 - \lambda_i t} \right) = - \lim_{u \rightarrow 0} \frac{d}{du} \prod_{i=1}^N (1 - \lambda_i t)^u$$

Overview of Derivation (cont.)

- Now consider the identity:

$$\prod_{i=1}^N (1 - \lambda_i t)^{-b} = {}_1\mathcal{F}_0^{(\alpha)} \left(b; t^N; \vec{\lambda} \right) = \sum_{\nu} \frac{[b]_{\nu}^{(\alpha)}}{|\nu|!} t^{|\nu|} C_{\nu}^{(\alpha)} \left(\vec{\lambda} \right)$$

- Our result follows from a substitution of the above identity.

$$\sum_{k \geq 1} \frac{p_k \left(\vec{\lambda} \right)}{k} t^k = - \lim_{u \rightarrow 0} \frac{d}{du} \left(\sum_{k \geq 1} \sum_{\nu \perp k} \frac{[-u]_{\nu}^{(\alpha)}}{|\nu|!} t^{|\nu|} C_{\nu}^{(\alpha)} \left(\vec{\lambda} \right) \right)$$

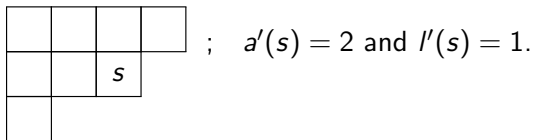
$$p_k \left(\vec{\lambda} \right) = \sum_{\nu \perp k} c_{\nu} C_{\nu}^{(\alpha)} \left(\vec{\lambda} \right)$$

The Coefficients c_{κ}

- The coefficients are:

$$\begin{aligned} c_{\nu} &= \frac{-1}{(|\nu| - 1)!} \left(\lim_{u \rightarrow 0} \frac{d}{du} [-u]_{\nu}^{(\alpha)} \right) \\ &= \frac{1}{(|\nu| - 1)!} \prod_{s \in \nu \setminus \{1,1\}} \left(a'(s) - \frac{l'(s)}{\alpha} \right) \end{aligned}$$

- For a box s in the Young diagram of ν , $a'(s)$ is the co-arm length and $l'(s)$ is the co-leg length e.g.



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